

From Eq.(2.25), we can obtain the Fourier transform of ψ as

$$\psi = \begin{cases} f^+(k_x, k_y) e^{jk_y y} & (y \geq 0) \\ f^-(k_x, k_y) e^{jk_y y} & (y \leq 0) \end{cases} \quad (2.26)$$

Then, in order to satisfy the radiation condition at $y \rightarrow \pm\infty$, we must choose

$$k_y^- = -k_y^+ = \begin{cases} j\sqrt{k_x^2 + k_y^2 - k^2} & (k \leq \sqrt{k_x^2 + k_y^2}) \\ -\sqrt{k^2 - k_x^2 - k_y^2} & (k \geq \sqrt{k_x^2 + k_y^2}) \end{cases} \quad (2.27)$$

As in the above aperture radiation example, we make the Fourier transform of the field components in Eq.(2.19) after setting $A_{z^-} = 0$ and $A_z = \psi$ and obtain

$$\hat{H}_x^+ = \frac{jk_x^2}{\mu_0} \psi^+, \quad \hat{E}_z^+ = -\frac{k_x k_y}{j\omega\mu_0 \epsilon_0} \psi^+, \quad (2.28a)$$

$$\hat{H}_x^- = -\frac{jk_x^2}{\mu_0} \psi^-, \quad \hat{E}_z^- = -\frac{k_x^2 k_y}{j\omega\mu_0 \epsilon_0} \psi^-, \quad (2.28b)$$

$$\hat{H}_y^+ = 0, \quad \hat{E}_z^+ = \frac{1}{j\omega\mu_0 \epsilon_0} [k^2 - k_x^2] \psi^+. \quad (2.28c)$$

According to the continuity of the tangential electric fields, i.e.,

$$E_x^+|_{y=0} = E_x^-|_{y=0} \quad \text{and} \quad E_z^+|_{y=0} = E_z^-|_{y=0}, \quad \text{we have}$$

$$f^+(k_x, k_y) = f^-(k_x, k_y) = f(k_x, k_y). \quad (2.29)$$

Also, according to the boundary conditions of the magnetic fields at $y=0$, i.e., $H_x^+ - H_x^- = -J_x$ or $\hat{H}_x^+ - \hat{H}_x^- = -\hat{J}_x$ when $y=0$, by using Eq.(2.26) and considering Eqs.(2.25) and (2.27), $f(k_x, k_y)$ can be determined by

$$f(k_x, k_y) = \frac{j\mu_0}{2k_x^+} \hat{J}_x, \quad (2.30)$$

Advanced Electromagnetic Theory

J Elliott



Advanced Electromagnetic Theory:

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